

Low-Loss Frequency-Agile Bandpass Filters With Controllable Bandwidth and Suppressed Second Harmonic

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Abstract—This paper presents a novel approach to design frequency-agile bandpass filters with constant absolute bandwidth and passband shape, as well as a suppressed second harmonic. A novel mixed electric and magnetic coupling scheme is proposed to control the coupling coefficient variation. Theoretical analysis indicates that it is able to achieve desired coupling coefficients between the proposed resonators at various frequencies so as to obtain constant absolute bandwidth. Moreover, this half-wavelength resonator has a Q higher than the quarter- and half-wavelength counterparts, thus resulting in low insertion loss. A filter of this type is designed to validate the proposed idea. To remove the spurious responses of the filter, a method is then introduced to suppress the second harmonic without degrading the passband performance. For demonstration, two frequency-agile filters with 60- and 80-MHz constant absolute bandwidth are implemented with the frequency tuning range from 680 to 1000 MHz. Comparisons of experimental and simulated results are presented to verify the theoretical predications.

Index Terms—Bandpass filter, constant bandwidth, harmonic suppression, Q factor, tunable filter.

I. INTRODUCTION

FREQUENCY-AGILE bandpass filters are essential for multiband and wideband systems. Extensive research has been done and different tuning devices have been employed. Semiconductor varactors are widely used due to the high tuning speed and reliability [1]–[11]. Unfortunately, the Q of varactor diodes is usually low, leading to high insertion loss. In contrast, microelectromechanical system (MEMS) varactors possess higher Q and linearity [12], [13]; these, however, are expensive. Ferroelectric components have recently attracted

much attention in tunable filter designs due to the tunability of the dielectric constant [14]–[16].

Regardless of the tuning devices used, it is desirable in certain applications to maintain constant absolute bandwidth and passband shape as the passband is tuned. Various methods have been utilized to achieve constant absolute bandwidth. Hunter and Rhodes [1] used combline filter topology realized in a suspended substrate stripline with the resonator's electrical length of 53° to obtain constant absolute bandwidth. This method can also be used in microstrip combline filters using stepped-impedance resonators [3]. Another method to achieve constant absolute bandwidth is to introduce a fixed or variable attenuation pole near the passband to force nearly constant bandwidth [4], [17]. The inter-stage coupling variation can also be controlled to obtain constant absolute bandwidth. In [5], a variable capacitor is utilized to control the coupling. In [6], an independent electric and magnetic coupling scheme is used to manipulate the coupling coefficient variation so that the bandwidth can be controlled.

Insertion loss is another important issue in frequency-agile bandpass filter designs. For a filter with a specific bandwidth, the insertion loss is mainly determined by the resonator Q . In order to reduce the insertion loss, high- Q components, e.g., MEMS components, are necessitated [12], [13], thus leading to high cost.

In this paper, a novel approach is proposed to design frequency-agile bandpass filters with constant absolute bandwidth and passband shape together with suppressed second harmonic. The proposed resonator is composed of an open-ended transmission line with one varactor loaded at one end. The variation of the inter-stage coupling coefficient can be controlled by using a novel mixed electric and magnetic coupling scheme; this allows that the requirement of constant absolute bandwidth can be met. Moreover, this half-wavelength resonator has higher overall Q than the widely used quarter- and half-wavelength counterparts. In turn, low-cost silicon varactors can be used to obtain acceptable performance. Using the proposed resonators and coupling scheme, a frequency-agile bandpass filters with 80-MHz constant absolute bandwidth is demonstrated with the frequency tuning range from 680 to 1000 MHz. A method is then introduced to suppress the second harmonic without affecting the passband performance. Two tunable filters with 60- and 80-MHz constant absolute bandwidth and suppressed second harmonic are implemented. Experimental and simulated results are presented to verify the proposed method.

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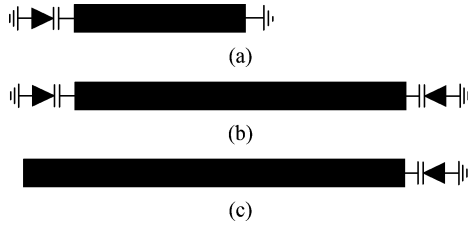


Fig. 1. (a) Conventional $\lambda/4$ resonator with one varactor. (b) Conventional $\lambda/2$ resonator with two varactors. (c) Proposed asymmetric $\lambda/2$ resonator with one varactor.

II. DESIGN THEORY

To design low-loss tunable bandpass filters with constant absolute bandwidth, the requirement on the unloaded resonator Q , coupling coefficients, and external Q should be satisfied.

A. Unloaded Resonator Q

It is common that tunable filter designs employ $\lambda/4$ and $\lambda/2$ resonators. Fig. 1(a) shows the configuration of a conventional $\lambda/4$ resonator. One end of the microstrip line is loaded by a varactor and the other end is shorted [2], [3], [6], [8]. Fig. 1(b) shows a conventional $\lambda/2$ resonator with two varactors loaded at the two ends, forming a symmetric structure with respect to the center plane of the transmission line [4], [10], [11]. These two types of resonators are widely used in tunable filter designs.

In this design, a novel $\lambda/2$ resonator is proposed, as shown in Fig. 1(c). Only one varactor is loaded at one end and the structure is asymmetric. The proposed resonator has higher Q than the other two counterparts. For comparative purposes, it is assumed that the Q values of all of the varactors are identical and lower than that of conventional microstrip line; the latter is true for most silicon varactors. The three resonators have the same fundamental resonant frequency under the same bias voltage. In this case, the microstrip line of the asymmetric $\lambda/2$ resonator is longer than the symmetric one. Since the microstrip has higher Q than the varactor, the overall Q of the proposed resonator is higher than that of the symmetric $\lambda/2$ resonator. On the other hand, the symmetric $\lambda/2$ resonator is equivalent to the $\lambda/4$ resonator at the fundamental resonant frequency. Hence, they have the same unloaded Q . As a result, the proposed resonator has higher unloaded Q than the others. For demonstration, a simulation study is conducted to extract the unloaded Q 's. The same varactor model and microstrip lines are utilized in the simulation and the three resonators are tuned to resonate at the same frequency of 850 MHz. The resonators in Fig. 1(a)–(c) have the unloaded Q 's of 58, 57, and 68, respectively. Hence, the proposed resonator has a Q higher than others. In turn, for given filter specifications, use of the proposed resonator can alleviate the requirement on the varactor Q , resulting in low-cost design.

B. Coupling Coefficient

For a second-order filter, the required coupling coefficient k is given by

$$k = \frac{BW}{f_0 \sqrt{g_1 g_2}} \quad (1)$$

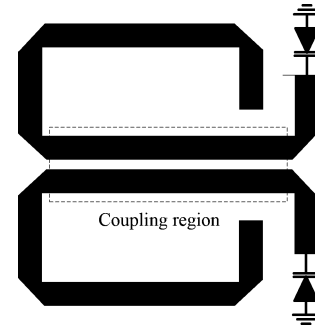


Fig. 2. Inter-stage coupling structure.

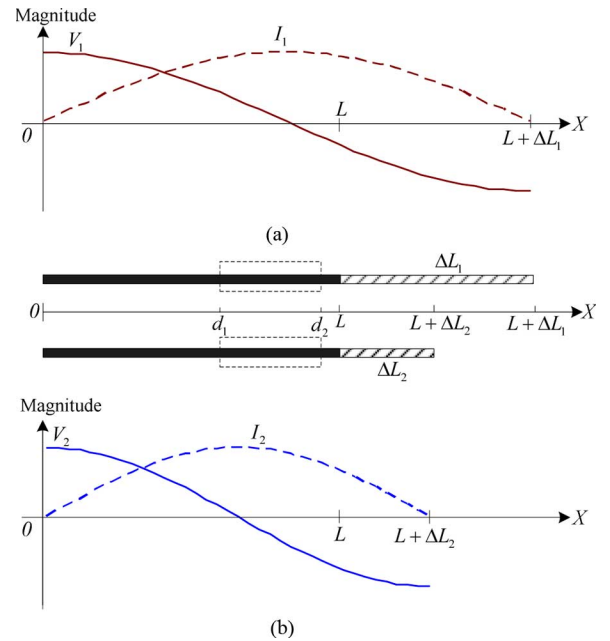


Fig. 3. Normalized voltage and current distributions on the resonator. (a) At lower frequency f_L . (b) At upper frequency f_U .

where BW represents the absolute bandwidth, g_1 and g_2 are the element values of the low-pass prototype filter, and f_0 is the center frequency. As can be seen, to achieve constant absolute bandwidth over the tuning range, the coupling coefficient should become lower as the center frequency increases.

To obtain the desired coupling coefficients, a novel mixed electric and magnetic coupling scheme is proposed, as shown in Fig. 2. For coupled microstrip lines, the electric and magnetic coupling can be evaluated by studying the voltage and current distributions [19]. Ignoring the parasitic effects, the varactor can be made equivalent to a segment of open-ended microstrip line. Fig. 3 shows the equivalent circuits of one resonator under different varactor bias voltages. The darkened parts denote real microstrip lines with the length of L . The filled parts (ΔL_1 and ΔL_2) are equivalent to the varactor under different bias voltages. The region between d_1 and d_2 represent the lines in the coupling region.

When the capacitance of the varactor is larger, the equivalent microstrip line is longer (ΔL_1) and the resonant frequency is

lower (f_L). At fundamental resonance, the normalized voltage and current can be expressed as

$$V_1(x) = \cos \beta_1 x \quad (2)$$

$$I_1(x) = \sin \beta_1 x \quad (3)$$

where β_1 is the propagation constant at f_L . The voltage and current distributions are illustrated in Fig. 3(a). In the case of smaller varactor capacitance, the equivalent microstrip line is shorter (ΔL_2) and the resonant frequency is higher (f_U). Similarly, the voltage and current can be given as

$$V_2(x) = \cos \beta_2 x \quad (4)$$

$$I_2(x) = \sin \beta_2 x \quad (5)$$

where β_2 is the propagation constant at f_U . The voltage and current distributions are shown in Fig. 3(b). The electric and magnetic coupling coefficients k_e and k_m in the two cases can be expressed as [19]

$$|k_{e,i}| = p \times \int_{d_1}^{d_2} |V_i(x)|^2 dx = p \times \int_{d_1}^{d_2} (\cos \beta_i x)^2 dx \quad (6)$$

$$|k_{m,i}| = p \times \int_{d_1}^{d_2} |I_i(x)|^2 dx = p \times \int_{d_1}^{d_2} (\sin \beta_i x)^2 dx \quad (7)$$

where p is a positive constant and $i = 1, 2$, corresponding to the lower and upper frequencies, respectively.

The overall coupling coefficient can vary inversely with frequency if a proper coupling region is selected. Since the electric and magnetic coupling coefficients are out of phase, when the magnetic coupling is dominant, the magnitude of coupling coefficient is

$$|k_i| = |k_{m,i}| - |k_{e,i}|. \quad (8)$$

When the magnetic coupling varies inversely with frequency while the electric coupling varies directly with frequency, the net coupling is proportional to frequency. Fig. 4 illustrates such a case. Inspecting (6) and (7), we can find that the coupling strength can be indicated by the area below the curve of $|V_i(x)|^2$ and $|I_i(x)|^2$ between d_1 and d_2 . From Fig. 4(a), it is observed that the electric coupling coefficient at lower frequency (indicated by the area below the dark red line (in online version) with cross symbols between d_1 and d_2) is lower than that at upper frequency (indicated by the area under the blue line (in online version) with square symbols between d_1 and d_2), namely,

$$|k_{e,1}| < |k_{e,2}|. \quad (9)$$

Similarly, Fig. 4(b) indicates that

$$|k_{m,1}| > |k_{m,2}|. \quad (10)$$

Inspecting (8)–(10) gives

$$|k_1| > |k_2|. \quad (11)$$

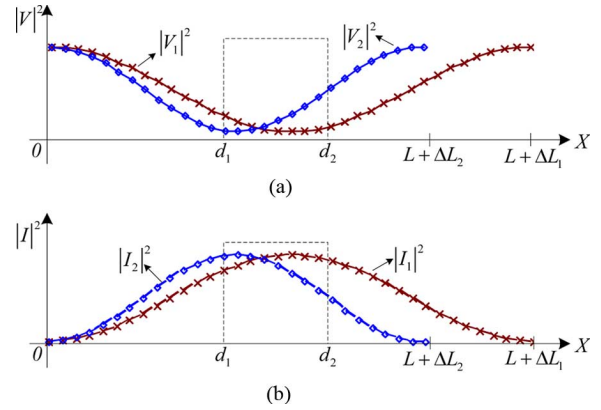


Fig. 4. Comparison of electric and magnetic coupling at various frequencies. (a) Electric coupling. (b) Magnetic coupling.

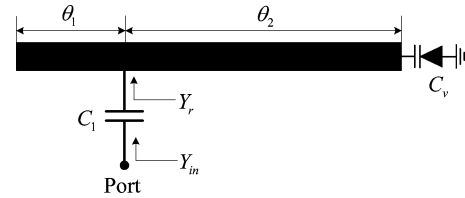


Fig. 5. Resonator with the input coupling network.

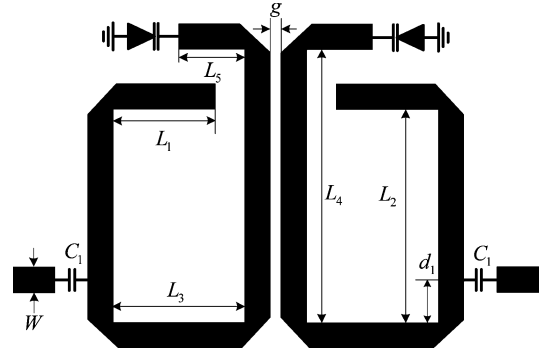


Fig. 6. Configuration of the tunable filter with constant absolute bandwidth.

Therefore, the coupling coefficient becomes lower as the center frequency is tuned upward; this can meet the requirement of constant absolute bandwidth.

C. External Q

The required external Q or Q_e for a second-order filter is given by

$$Q_e = \frac{f_0}{\text{BW} \sqrt{g_0 g_1}}. \quad (12)$$

To maintain constant absolute bandwidth, Q_e should increase as the frequency shifts upward. This can be realized by using the structure shown in Fig. 5.

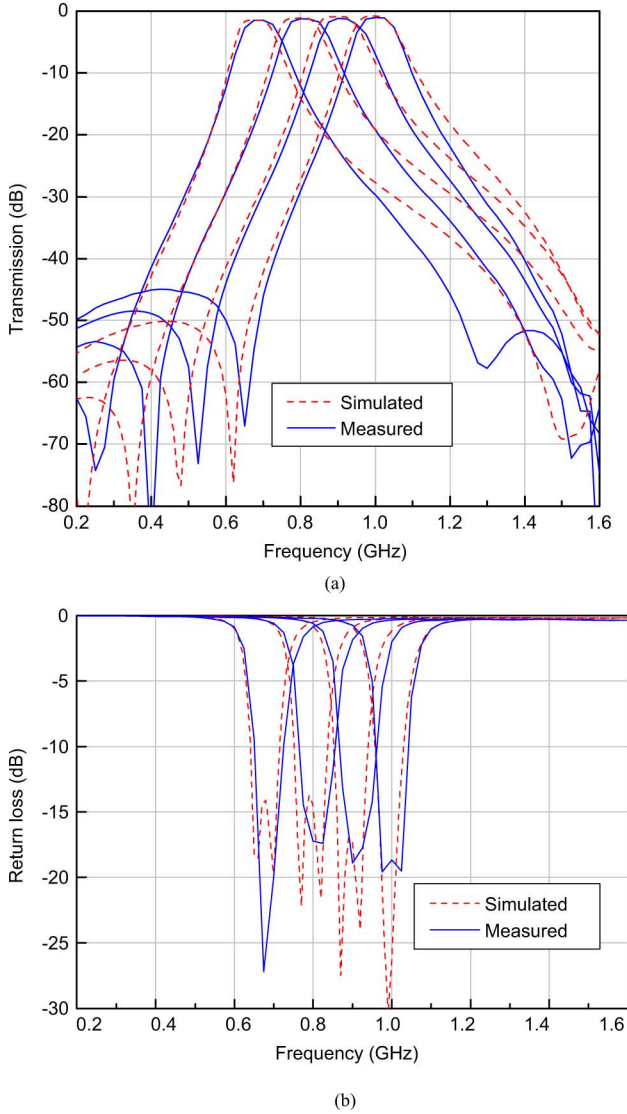


Fig. 7. Simulated and measured responses. (a) Transmission. (b) Return loss.

To facilitate the analysis, the parasitic effects of the varactor and the line discontinuity are ignored. The input admittance of the resonator Y_r , seen from the input port, is

$$Y_r = \frac{jY}{Y - \omega C_v \tan \theta_2} [Y(\tan \theta_1 + \tan \theta_2) + \omega C_v (1 - \tan \theta_1 \tan \theta_2)] \quad (13)$$

where Y is the characteristic admittance of the microstrip line, θ_1 and θ_2 are the electric length, and C_v is the capacitance of the varactor. The overall input admittance Y_{in} is

$$Y_{in} = \frac{j\omega C_1 Y_r}{j\omega C_1 + Y_r} \quad (14)$$

Q_e is then determined as

$$Q_e = \frac{\omega_0}{2Y_0} \left. \frac{\partial \text{Im}[Y_{in}]}{\partial \omega} \right|_{\omega=\omega_0} \quad (15)$$

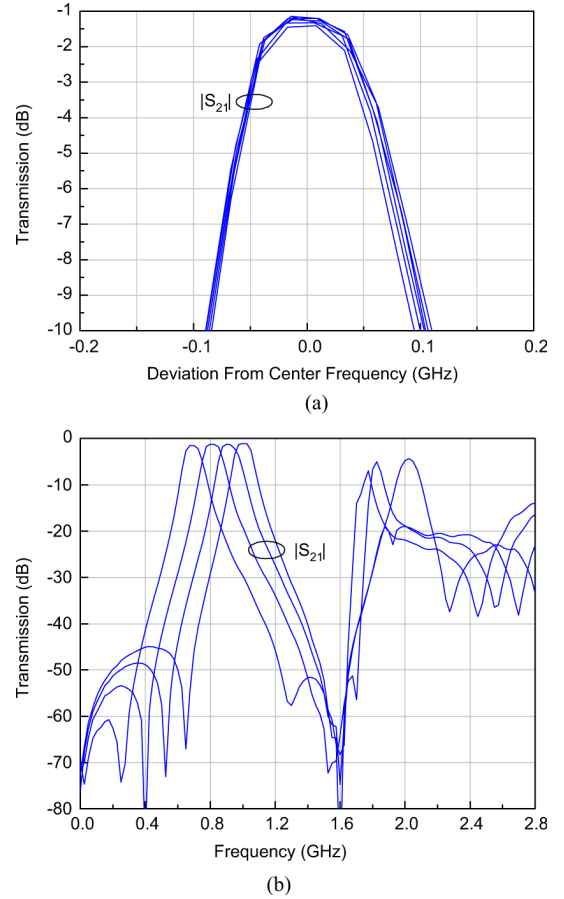


Fig. 8. Measured responses. (a) Superposition of the passband responses. (b) Wideband responses.

TABLE I
COMPARISON WITH OTHER DESIGN

	Resonator	Varactor Q	Loss tangent	Unloaded Q
[6]	$\lambda/4$ with one varactor	Around 2000	0.0009	53-152
This work	$\lambda/2$ with one varactor	35-100	0.0018	42-72

It can be seen that the capacitance C_1 and the electrical length θ_1 and θ_2 will affect Q_e . By properly choosing these three parameters, the desired Q_e can be obtained within the frequency tuning range.

III. FREQUENCY-AGILE FILTER IMPLEMENTATION

Based on the previous design theory, a frequency-agile filter with constant absolute bandwidth is implemented. The filter configuration is shown in Fig. 6. This filter employs two resonators. The silicon varactors are 1sv277 from Toshiba, Tokyo, Japan. The input and output feeding lines are tapped at the resonators through the capacitors. The tap position d_1 and the capacitance C_1 are tuned so as to fulfill the requirement of Q_e over the frequency tuning range.

The filter is fabricated on the substrate with the thickness of 0.82 mm, relative dielectric constant of 6.03, and loss tangent of 0.0018. The dimensions are determined as follows: $L_1 = 5.0$ mm, $L_2 = 18.8$ mm, $L_3 = 6.8$ mm, $L_4 = 22.0$ mm,

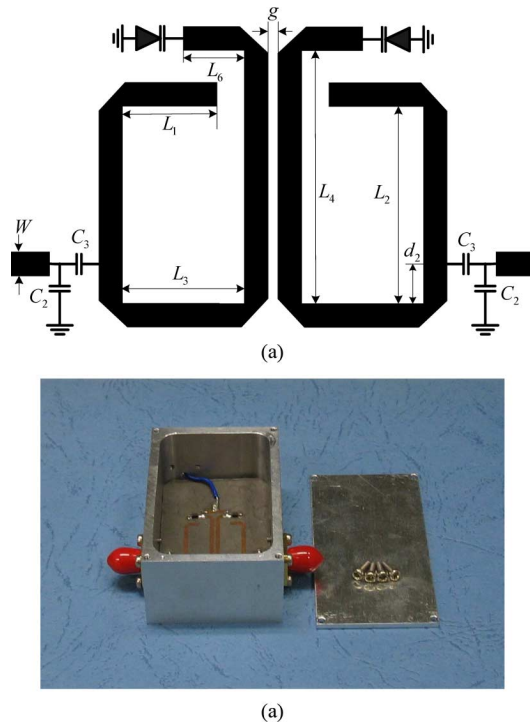


Fig. 9. Tunable filter with the suppressed second harmonic responses. (a) Filter configuration. (b) Photograph of the fabricated filter.

$L_5 = 3.5$ mm, $W = 1.2$ mm, $d_1 = 2.5$ mm, $g = 0.5$ mm, and $C_1 = 3$ pF. The filter is enclosed in a metal cavity with the dimensions of $5 \times 3.4 \times 2$ cm³.

Fig. 7 shows the simulated and measured responses, which agree reasonably well with each other. As shown in Fig. 7(b), the center frequency can be tuned between 680 and 1000 MHz, featuring the fractional tuning range of 38%. The measured return loss over the tuning range is greater than 15 dB. For each tuning state, there are two transmission zeros near the passband edges, which improve the roll-off rate. These transmission zeros are due to the tap connections of input and output port, which lead to quarter-wavelength resonance. Fig. 8(a) shows a superposition of the measured passband responses at various frequencies. As can be observed, the passband shape and absolute bandwidth are kept nearly constant over the tuning range. The measured 1-dB absolute bandwidth is 80 ± 3.5 MHz. The variation is less than 4.4%, which is lower than those (7%–10%) in [1] and [6] and a little bit higher than those (3%–3.5%) in [3] and [4]. Therefore, the bandwidth can be considered as constant. For all the tuning states, the insertion loss is from 1.1 to 1.5 dB, thus indicating low insertion loss. The variation of insertion-loss in absolute value is 4.6%, which is lower than the value of 7.2% in [5] and 11.4% in [6].

This design is compared with that in [6], as tabulated in Table I. The measured unloaded Q of this filter is 42–72 from 680 to 1000 MHz. This is somewhat lower than the reported one of 53–152 from 850 to 1400 MHz in [6]. However, the design in [6] employed high- Q GaAs varactors and low-loss substrate. In contrast, this is a low-cost design, which utilizes ordinary silicon varactors and substrate. This indicates the superiority of the proposed resonator.

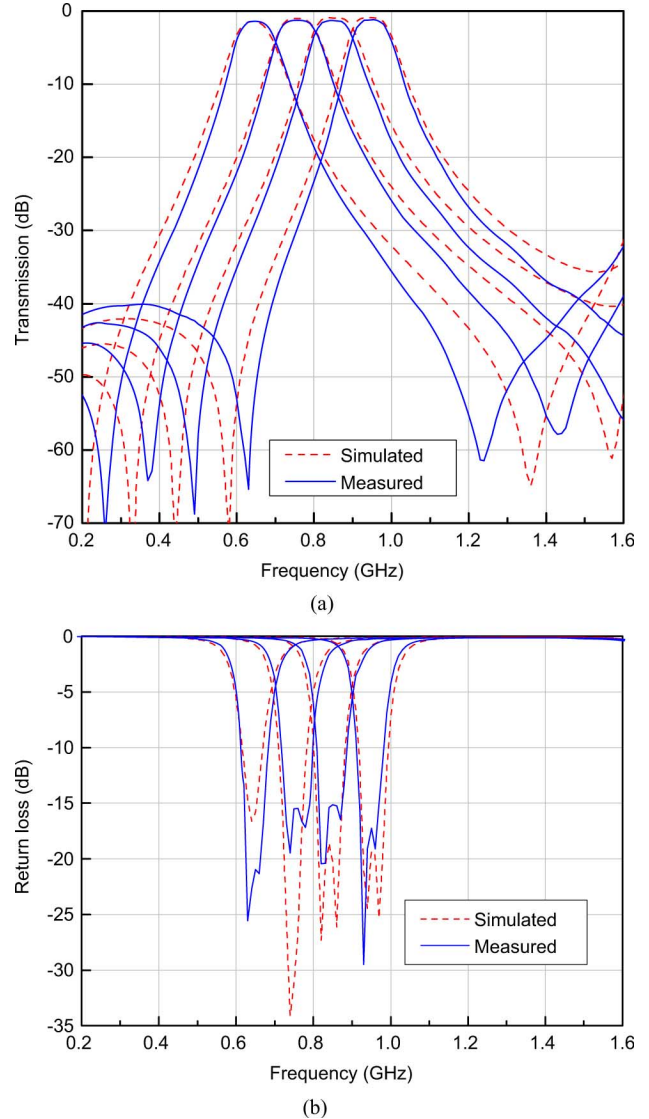


Fig. 10. Responses of the filter with 80-MHz bandwidth. (a) Transmission. (b) Return loss.

Despite the good passband performance of the implemented filter, it suffers from high second harmonic levels, as shown in Fig. 8(b). It is straightforward to suppress them using a low-pass or bandstop structure. However, it will induce extra insertion loss and increase the circuit size. In applications, it is desirable to reduce the harmonic levels without affecting the passband performance. This leads to the following designs.

IV. FREQUENCY-AGILE FILTER WITH SUPPRESSED HARMONIC

The second harmonic of the previous filter can be suppressed by using input and output coupling networks with the intrinsic bandpass characteristic. Fig. 9 shows such a filter configuration. This filter is the same as the previous one aside from the input and output coupling networks. The shunt capacitor C_2 and series capacitor C_3 form a bandpass network. When the second harmonic is within the stopband of this bandpass network, they can be suppressed. On the other hand, the coupling networks affect Q_e . Based on (14), the input admittance can be derived as

$$Y'_{in} = \frac{j\omega Y_r (C_2 + C_3) - \omega^2 C_2 C_3}{j\omega C_3 + Y_r}. \quad (16)$$

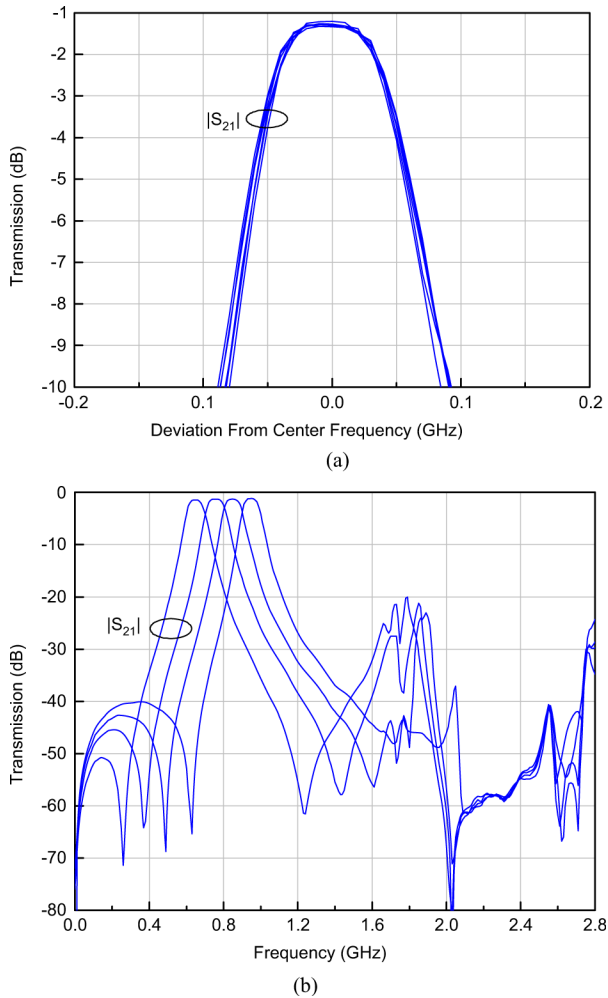


Fig. 11. Measured responses of the filter with 80-MHz bandwidth. (a) Superposition of the passband responses. (b) Wideband responses.

Inspecting (15) and (16), we can find that the two capacitors will affect Q_e . Therefore, the capacitors C_2 and C_3 should fulfill the requirements of both Q_e and harmonic suppression. It is noted that Q_e is controlled not only by the capacitance C_2 and C_3 , but also by the tap position d_2 . Hence, there are sufficient degrees of freedom to meet the requirements of both Q_e and harmonic suppression. In this manner, the harmonic can be suppressed without degrading passband performance.

Using the configuration, two filters with 80- and 60-MHz absolute bandwidth are implemented with the suppressed second harmonic. The two filters are fabricated on the substrate with the relative dielectric constant of 6.03, thickness of 0.82 mm, and loss tangent of 0.0018. The filter with 80-MHz bandwidth has the following dimensions: $L_1 = 5.0$ mm, $L_2 = 18.8$ mm, $L_3 = 6.8$ mm, $L_4 = 22.0$ mm, $L_6 = 5.0$ mm, $W = 1.2$ mm, $d_2 = 2.8$ mm, $g = 0.5$ mm, $C_2 = 3.3$ pF, and $C_3 = 6.8$ pF. The overall size is around $0.12 \lambda_g \times 0.12 \lambda_g$, where λ_g is the guided wavelength at lowest passband frequency. The filter is enclosed in a metal cavity with the dimensions of $5 \times 3.4 \times 2$ cm³. The fabricated filter is shown in Fig. 9(b).

Fig. 10 illustrates the simulated and measured results. The passband tuning range is from 650 to 960 MHz, featuring a 38% fractional tuning range. The measured in-band return loss

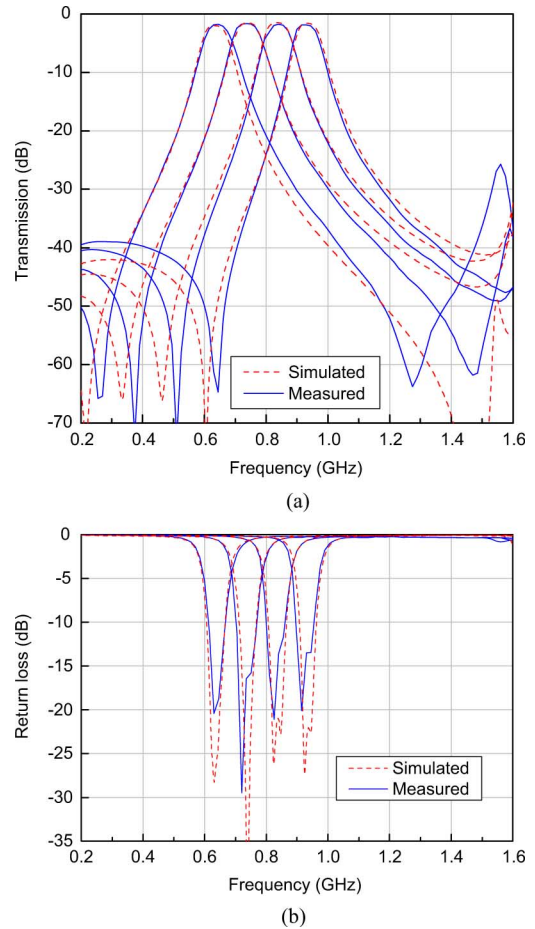


Fig. 12. Responses of the filter with 60-MHz bandwidth. (a) Transmission. (b) Return loss.

is greater than 15 dB for all the tuning states. The measured 1-dB absolute bandwidth is 80 ± 3.5 MHz. The passband shape and insertion loss are maintained nearly constant over the tuning range, as shown in Fig. 11(a). The insertion loss ranges from 1.2 to 1.5 dB, which is nearly the same as the previous one; this indicates that the passband performance is nearly not affected. Fig. 11(b) shows the measured wideband responses. It can be seen that the second harmonic levels are below -20 dB. Hence, it can be summarized that the passband performance is preserved, while the second harmonic is suppressed. The input third-order intermodulation intercept point (IIP3) is measured at 2.5-V bias with 1-MHz frequency spacing. The IIP3 is measured at 13 dBm.

To demonstrate the wide applicability of this method, a filter with 60-MHz bandwidth is also implemented. The dimensions are as follows: $L_1 = 3.5$ mm, $L_2 = 16.8$ mm, $L_3 = 6.8$ mm, $L_4 = 20$ mm, $L_6 = 5.0$ mm, $W = 1.2$ mm, $d_2 = 3.5$ mm, $g = 0.55$ mm, $C_2 = 5$ pF, and $C_3 = 8$ pF. The filter is also enclosed in a metal cavity with the dimensions of $5 \times 3.4 \times 2$ cm³. Fig. 12 shows the simulated and measured results. The passband frequency can be tuned from 630 to 930 MHz, featuring the fractional tuning range of 38.5%. The measured in-band return loss is around 20 dB for all the tuning states. The 1-dB absolute bandwidth is 60 ± 3 MHz. The passband shape and insertion loss are maintained nearly constant over the tuning

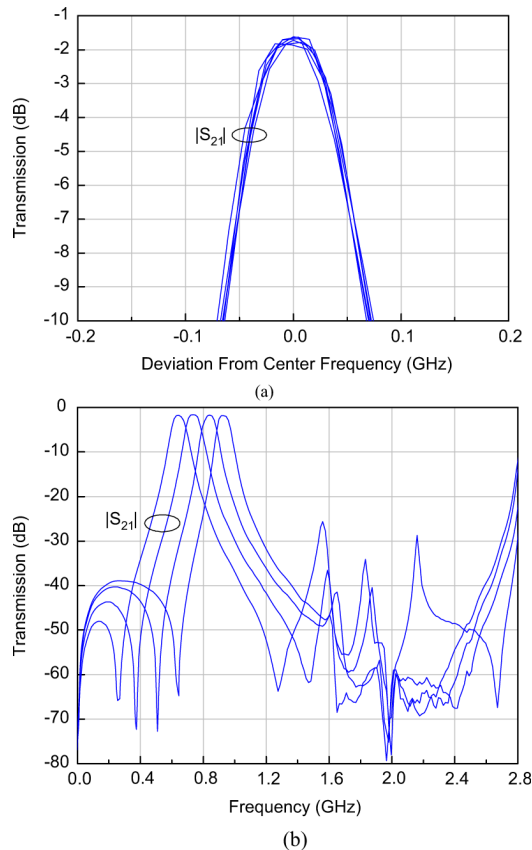


Fig. 13. Measured responses of the filter with 60-MHz bandwidth. (a) Superposition of the passband responses. (b) Wideband responses.

range, as shown in Fig. 13(a). The insertion loss ranges from 1.6 to 2.0 dB. Fig. 13(b) shows the measured wideband responses. It can be seen that the second harmonic levels are suppressed below -25 dB. The measured IIP3 is also around 13 dBm.

From the two design examples, it is observed that the second harmonic can be suppressed by simple and compact circuits and the passband performance is not affected. The absolute bandwidth and passband shape can be maintained constant over the tuning range. Moreover, the bandwidth can be controlled by changing the design parameters, indicating this topology can be used to achieve various bandwidth specifications.

V. CONCLUSION

In this paper, a novel method has been presented to design frequency-agile bandpass filters with constant absolute bandwidth and passband shape together with the suppressed second harmonic. Both theory and experiments have been provided, showing that Q_e and coupling coefficient variation can be controlled to fulfill the requirement of constant absolute bandwidth. The proposed resonator has a Q higher than the quarter- and half-wavelength counterparts, which, in turn, can result in low insertion loss. With the intrinsic bandpass characteristic, the compact input and output coupling networks are utilized to not only obtain desired Q_e variation, but also reduce the second harmonic. The experimental results have been presented, showing that constant absolute bandwidth and passband shape can be

achieved and the second harmonic can be suppressed without degrading the passband performance. Furthermore, the absolute bandwidth can be controlled by altering the design parameters, indicating the wide applicability of this method. Transmission zeros have been generated near the passband edges, resulting in high skirt selectivity. With these features, this kind of tunable filter will be useful in multiband and wideband systems.

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